

Published by NIGERIAN SOCIETY OF PHYSICAL SCIENCES Available online @ https://journal.nsps.org.ng/index.php/jnsps

J. Nig. Soc. Phys. Sci. 5 (2023) 1484

A class of single-step hybrid block methods with equally spaced points for general third-order ordinary differential equations

Maduabuchi Gabriel Orakwelu^{a,b,*}, Olumuyiwa Otegbeye^c, Hermane Mambili-Mamboundou^a

^aSchool of Mathematics, Statistics and Computer Science, University of KwaZulu-Natal, Pietermaritzburg, South Africa ^bDSI-NRF Centre of Excellence in Mathematical and Statistical Sciences (CoE-MaSS), South Africa.

^cSchool of Computer Science and Applied Mathematics, Faculty of Science, University of the Witwatersrand, Johannesburg, South Africa

Abstract

This study presents a class of single-step, self-starting hybrid block methods for directly solving general third-order ordinary differential equations (ODEs) without reduction to first order equations. The methods are developed through interpolation and collocation at systematically selected evenly spaced nodes with the aim of boosting the accuracy of the methods. The zero stability, consistency and convergence of the algorithms are established. Scalar and systems of linear and nonlinear ODEs are approximated to test the effectiveness of the schemes, and the results obtained are compared against other methods from the literature. Significantly, the study shows that an increase in the number of intra-step points improves the accuracy of the solutions obtained using the proposed methods.

DOI:10.46481/jnsps.2023.1484

Keywords: Third-order problems, Convergence, Hybrid block methods, Boundary value problems, Initial value problems

Article History : Received: 01 April 2023 Received in revised form: 08 October 2023 Accepted for publication: 16 October 2023 Published: 28 November 2023

© 2023 The Author(s). Published by the Nigerian Society of Physical Sciences under the terms of the Creative Commons Attribution 4.0 International license. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Communicated by: O. Adeyeye

1. Introduction

Many problems in the fields of mathematics and engineering are formulated using initial and boundary value problems (IVP and BVP) for third-order ODEs. Most physical problems, including the deflection of a curved beam with a constant or varying cross-section, a three-layer beam, the motion of a rocket, thin film flow, electromagnetic waves, gravity-driven flows, physical oceanography, and the context of a variational inequality involve the use of third-order ODEs of the form

$$y'''(x) = f(x, y, y', y''),$$
(1)

subject to initial conditions

$$y(x_0) = y_0, y'(x_0) = \delta_0, y''(x_0) = \omega_0,$$

boundary conditions

$$y(x_0) = y_0, y'(x_0) = \delta_0, y'(x_N) = \omega_0,$$

or with mixed boundary conditions

$$y(x_0) = y_0, y'(x_0) = \delta_0, y(x_N) - y'(x_N) = \omega_0,$$

where $x_0, x_N, y_0, \delta_0, \omega_0 \in \mathfrak{R}$, and *f* is a continuous function and satisfies a Lipschitz condition as given Henrici [1].

Researchers have published a considerable amount of literature on Linear Multistep Methods (LMM) for solving Equation

^{*}Corresponding author tel. no: +2348034509462

Email address: maduabuchi@gmail.com (Maduabuchi Gabriel Orakwelu)

(1). Through Taylor expansions, some 3 and 5-step schemes for the special third-order ODEs were presented by Rajabi *et al.* [2]. A one-step block method with four equidistant generalized hybrid points was presented by Adeyeye and Omar [3]. Via numerical integration, a two-point four-step direct implicit block method implemented at two points simultaneously in a block using four backward steps was proposed by Majid *et al.* [4]. Also, a fourth and fifth derivative, three-point implicit block method was presented by Allogmany and Ismail [5]. Using interpolation and collocation techniques, implicit continuous LMM for solving Equation (1) was presented by Jator [6]. Linear multistep methods have also been used to solve other types of differential equations in Refs. [7, 8].

Proposed to overcome the Dahlquist Barrier theorem, hybrid methods were introduced and have continued to generate interest among numerical analysts as given in Ref. [9, 10]. In this paper, we derive a class of hybrid block methods via interpolation and collocation technique, that directly solves general third-order ODEs. We reiterate that the most typical methods for attempting to solve (1) typically entail reducing the problem to a system of first-order differential equations and then solving the system using one of the available methods, which has been shown to have some significant computational drawbacks like requiring more time and labour from the user as as presented by Jator *et al.* [11]. Within the decade researchers have clearly shown an unequivocal dependence of the accuracy of block hybrid methods on the number and type of grid points incorporated in the derivation process as presented in Refs. [12-15].

In this paper, a self-starting class of single-step hybrid block methods are presented for numerically integrating general thirdorder ODEs. We formulate certain members of the single-step hybrid block methods, prove the convergence and perform numerical tests to check the precision of the suggested methods by contrasting them with other schemes from the literature.

2. Derivation and Analysis of Method

2.1. Derivation of the Methods

In this section, we describe the formulation of a class of onestep block hybrid methods via the interpolation and collocation approach as proposed by Onumanyi *et al.* [16], which will be utilized to generate a number of discrete solutions for solving (1).

We start by deriving a block hybrid method of the form

$$Y(x) = \alpha_0(x)y_n + \alpha_{\frac{1}{2}}(x)y_{n+\frac{1}{2}} + \alpha_1(x)y_{n+1} + h^3 \left(\sum_{j=0}^1 \beta_j(x)y_{n+j} + \sum_{\nu=1}^m \beta_{p_\nu}(x)f_{n+p_\nu}\right), \quad (2)$$

where $p_v \in (0, k)$ are a countable number $m \in \mathbb{O}$ of equally spaced, non-integer off-grid points with $\frac{1}{2}$ as midpoint derived by

$$x_{n+p_{v}} = x_{n} + \frac{v}{m+1}h, v = 1, 2, \cdots, m.$$
 (3)

Equation (1) contains the first and second derivative, hence, the derivatives of Equation (2) are given as

$$Y'(x) = \frac{1}{h} \left(\alpha'_0(x) y_n + \alpha'_{\frac{1}{2}}(x) y_{n+\frac{1}{2}} + \alpha'_1(x) y_{n+1} + h^3 \left(\sum_{j=0}^{1} \beta'_j(x) y_{n+j} + \sum_{\nu=1}^{m} \beta'_{p_\nu}(x) f_{n+p_\nu} \right) \right), \quad (4)$$

$$Y^{''}(x) = \frac{1}{h^2} \Big(\alpha_0^{\prime\prime}(x) y_n + \alpha_{\frac{1}{2}}^{\prime\prime}(x) y_{n+\frac{1}{2}} + \alpha_1^{\prime\prime}(x) y_{n+1} \\ + h^3 \Big(\sum_{j=0}^{1} \beta_j^{\prime\prime}(x) y_{n+j} + \sum_{\nu=1}^{m} \beta_{p_\nu}^{\prime\prime}(x) f_{n+p_\nu} \Big) \Big).$$
(5)

The following conditions are imposed on Equations (4) and (5)

$$Y'(x) = \delta(x), \quad Y''(x) = \omega(x), \tag{6}$$

$$Y'(a) = \delta_0, \quad Y''(a) = \omega_0.$$
 (7)

2.2. Specification of the methods

In this section, Equation (2) is utilized in obtaining a particular one-step block hybrid method with equally spaced off-grid points by specifying m. For illustrative purposes, the results for m = 3 is provided.

Evaluating Equation (2) at $x = \{x_{n+\frac{1}{4}}, x_{n+\frac{1}{2}}, x_{n+\frac{3}{4}}\}$, we generate the two main methods (8) and (9) given below as:

$$\begin{cases} \frac{h^{3}}{30720} \left(f_{n} + f_{n+1} + 126f_{n+\frac{1}{2}} + 86f_{n+\frac{1}{4}} + 26f_{n+\frac{3}{4}} \right) \\ + \frac{1}{8} \left(3y_{n} - y_{n+1} + 6y_{n+\frac{1}{2}} - 8y_{n+\frac{1}{4}} \right) \end{cases} = 0,(8)$$
$$\begin{cases} \frac{h^{3}}{30720} \left(-f_{n} - f_{n+1} - 126f_{n+\frac{1}{2}} - 26f_{n+\frac{1}{4}} - 86f_{n+\frac{3}{4}} \right) \\ + \frac{1}{8} \left(-y_{n} + 3y_{n+1} + 6y_{n+\frac{1}{2}} - 8y_{n+\frac{3}{4}} \right) \end{cases} = 0,(9)$$

The starters Equations (10) and (11) are obtained from Equation (7) and are given as:

$$\begin{cases} \frac{h^3}{10080} \left(-41f_n + f_{n+1} - 288f_{n+\frac{1}{2}} - 424f_{n+\frac{1}{4}} \right) \\ -88f_{n+\frac{3}{4}} + h\delta_n + 3y_n + y_{n+1} - 4y_{n+\frac{1}{2}} \end{cases} = 0, \quad (10)$$

$$\begin{cases} \frac{1}{720} h^3 \left(59f_n - 3f_{n+1} + 48f_{n+\frac{1}{2}} + 216f_{n+\frac{1}{4}} \right) \\ + 40f_{n+\frac{3}{4}} + h^2\omega_n - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \end{cases} = 0. \quad (11)$$

It is worth mentioning that the derivatives are generated by $\delta(x_{n+\tau_m}) = \delta_{n+\tau_m}$ and $\omega(x_{n+\tau_m}) = \omega_{n+\tau_m}$ as follows:

$$\begin{cases} \frac{h^3}{161280} \left(79f_n - 5f_{n+1} + 54f_{n+\frac{1}{2}} + 1532f_{n+\frac{1}{4}} \right) \\ + 20f_{n+\frac{3}{4}} + h\delta_{n+\frac{1}{4}} + 2\left(y_n - y_{n+\frac{1}{2}}\right) \end{cases} = 0, \quad (12)$$

$$\begin{cases} \frac{h^{3}}{10080} \left(f_{n} + f_{n+1} + 258f_{n+\frac{1}{2}} + 80f_{n+\frac{1}{4}} \right) = 0, \quad (13) \\ + 80f_{n+\frac{3}{4}} + h\delta_{n+\frac{1}{2}} + y_{n} - y_{n+1} \end{cases} = 0, \quad (13) \\ \begin{cases} \frac{h^{3}}{161280} \left(-5f_{n} + 79f_{n+1} + 54f_{n+\frac{1}{2}} + 20f_{n+\frac{1}{4}} \right) \\ + 1532f_{n+\frac{3}{4}} + h\delta_{n+\frac{3}{4}} + 2\left(y_{n+\frac{1}{2}} - y_{n+1} \right) \end{cases} = 0, \quad (14) \\ \begin{cases} \frac{h^{3}}{10080} \left(f_{n} - 41f_{n+1} - 288f_{n+\frac{1}{2}} - 88f_{n+\frac{1}{4}} \right) \\ - 424f_{n+\frac{3}{4}} + h\delta_{n+1} - y_{n} - 3y_{n+1} + 4y_{n+\frac{1}{2}} \end{cases} = 0, \quad (15) \end{cases} \\ \begin{cases} \frac{h^{3}}{2880} \left(-15f_{n} + 7f_{n+1} + 456f_{n+\frac{1}{2}} + 218f_{n+\frac{1}{4}} \right) \\ + 54f_{n+\frac{3}{4}} + h^{2}\omega_{n+\frac{1}{4}} - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{2}} \right) \end{cases} = 0, \quad (16) \end{cases} \\ \begin{cases} \frac{1}{720}h^{3} \left(f_{n} - f_{n+1} - 32f_{n+\frac{1}{4}} + 32f_{n+\frac{3}{4}} \right) \\ + h^{2}\omega_{n+\frac{1}{2}} - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{2}} \right) \end{cases} = 0, \quad (17) \end{cases} \\ \begin{cases} \frac{h^{3}}{2880} \left(-7f_{n} + 15f_{n+1} - 456f_{n+\frac{1}{2}} - 54f_{n+\frac{1}{4}} \right) \\ - 218f_{n+\frac{3}{4}} + h^{2}\omega_{n+\frac{3}{4}} - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{2}} \right) \end{cases} = 0, \quad (19) \end{cases} \end{cases}$$

For the one-step block hybrid method with m = 5 equally spaced off-grid points, we evaluate Equation (2) at

$$x = \left\{ x_{n+\frac{1}{6}}, \ x_{n+\frac{1}{3}}, \ x_{n+\frac{1}{2}}, \ x_{n+\frac{2}{3}}, \ x_{n+\frac{5}{6}}, \ x_{n+1} \right\}.$$

We generate the four main methods (20), (21),(22) and (23) given below as:

$$\begin{cases} \frac{h^{3}}{19595520} \left(412f_{n} + 97f_{n+1} + 44050f_{n+\frac{1}{2}} + 57525f_{n+\frac{1}{3}} + 20670f_{n+\frac{2}{3}} + 23673f_{n+\frac{1}{6}} + 57525f_{n+\frac{1}{3}} + 20670f_{n+\frac{2}{3}} + 23673f_{n+\frac{1}{6}} + 4773f_{n+\frac{5}{6}} \right) + \frac{1}{9} \left(5y_{n} - y_{n+1} + 5y_{n+\frac{1}{2}} - 9y_{n+\frac{1}{6}} \right) \\ \left\{ \frac{h^{3}}{19595520} (191f_{n} + 65f_{n+1} + 45824f_{n+\frac{1}{2}} + 40539f_{n+\frac{1}{3}} + 19749f_{n+\frac{2}{3}} + 9564f_{n+\frac{1}{6}} + 5028f_{n+\frac{5}{6}} \right) + \frac{1}{9} (2y_{n} - y_{n+1} + 8y_{n+\frac{1}{2}} - 9y_{n+\frac{1}{3}}) \\ \left\{ \frac{h^{3}}{19595520} \left(-65f_{n} - 191f_{n+1} - 45824f_{n+\frac{1}{2}} - 19749f_{n+\frac{1}{3}} - 40539f_{n+\frac{2}{3}} - 5028f_{n+\frac{1}{6}} - 9564f_{n+\frac{5}{6}} \right) + \frac{1}{9} \left(-y_{n} + 2y_{n+1} + 8y_{n+\frac{1}{2}} - 9y_{n+\frac{2}{3}} \right) \\ = 0, \quad (22)$$

$$\begin{cases} \frac{h^{3}}{19595520} \left(-97f_{n} - 412f_{n+1} - 44050f_{n+\frac{1}{2}}\right) \\ -20670f_{n+\frac{1}{3}} - 57525f_{n+\frac{2}{3}} - 4773f_{n+\frac{1}{6}} \\ -23673f_{n+\frac{5}{6}}\right) + \frac{1}{9} \left(-y_{n} + 5y_{n+1} + 5y_{n+\frac{1}{2}}\right) \\ -9y_{n+\frac{5}{6}}\right) \end{cases} = 0.$$
 (23)

The starters Equations (24) and (25) are obtained from Equation (7) and are given as:

$$\begin{cases} \frac{h^3}{268800} \left(-459f_n + 11f_{n+1} - 6400f_{n+\frac{1}{2}} - 6615f_{n+\frac{1}{3}} \right) \\ - 2025f_{n+\frac{2}{3}} - 6156f_{n+\frac{1}{6}} - 756f_{n+\frac{5}{6}} \right) + h\delta_n + 3y_n \\ + y_{n+1} - 4y_{n+\frac{1}{2}} \end{cases} = 0,(24)$$
$$\begin{cases} \frac{h^3}{26880} \left(1359f_n - 47f_{n+1} + 4352f_{n+\frac{1}{2}} + 1035f_{n+\frac{1}{3}} \right) \\ - 171f_{n+\frac{2}{3}} + 6300f_{n+\frac{1}{6}} + 612f_{n+\frac{5}{6}} \right) + h^2\omega_n \\ - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \end{cases} = 0.(25)$$

It is worth mentioning that the derivatives are generated by $\delta(x_{n+\tau_m}) = \delta_{n+\tau_m}$ and $\omega(x_{n+\tau_m}) = \omega_{n+\tau_m}$

$$\begin{cases} \frac{h^{3}}{21772800} \left(3541f_{n} - 479f_{n+1} - 143300f_{n+\frac{1}{2}} \right) \\ - 138825f_{n+\frac{1}{3}} - 71565f_{n+\frac{2}{3}} + 63114f_{n+\frac{1}{6}} \\ - 14886f_{n+\frac{5}{6}} \right) + h\delta_{n+\frac{1}{6}} + \frac{1}{3} \left(7y_{n} + y_{n+1} \right) \\ - 8y_{n+\frac{1}{2}} \right) \end{cases} = 0, (26)$$

$$\begin{cases} \frac{h^{3}}{21772800} \left(1111f_{n} + 361f_{n+1} + 152320f_{n+\frac{1}{2}} \right) \\ + 281235f_{n+\frac{1}{3}} + 69165f_{n+\frac{2}{3}} + 84924f_{n+\frac{1}{6}} \\ + 15684f_{n+\frac{5}{6}} \right) + h\delta_{n+\frac{1}{3}} + \frac{1}{3} \left(5y_{n} - y_{n+1} - 4y_{n+\frac{1}{2}} \right) \end{cases} = 0, (27)$$

$$\begin{cases} \frac{h^{3}}{268800} (11f_{n} + 11f_{n+1} + 4860f_{n+\frac{1}{2}} + 2565f_{n+\frac{1}{3}} \\ + 2565f_{n+\frac{2}{3}} + 594f_{n+\frac{1}{6}} + 594f_{n+\frac{5}{6}} \right) + h\delta_{n+\frac{1}{2}} + y_{n} \\ - y_{n+1} \end{cases} = 0, (28)$$

$$\begin{cases} \frac{h^{3}}{21772800} \left(361f_{n} + 1111f_{n+1} + 152320f_{n+\frac{1}{2}} \\ + 69165f_{n+\frac{1}{3}} + 281235f_{n+\frac{2}{3}} + 15684f_{n+\frac{1}{6}} \\ + 84924f_{n+\frac{5}{6}} \right) + h\delta_{n+\frac{2}{3}} + \frac{1}{3} \left(y_{n} - 5y_{n+1} + 4y_{n+\frac{1}{2}} \right) \end{cases} = 0, (29)$$

$$\begin{cases} \frac{h^{3}}{21772800} \left(-479f_{n} + 3541f_{n+1} - 143300f_{n+\frac{1}{2}} \\ -71565f_{n+\frac{1}{3}} - 138825f_{n+\frac{2}{3}} - 14886f_{n+\frac{1}{6}} \\ + 63114f_{n+\frac{5}{6}} \right) + h\delta_{n+\frac{5}{6}} + \frac{1}{3} \left(-y_{n} - 7y_{n+1} \\ + 8y_{n+\frac{1}{2}} \right) \end{cases} = 0, (30)$$

$$\begin{cases} \frac{h^3}{268800} (11f_n - 459f_{n+1} - 6400f_{n+\frac{1}{2}} - 2025f_{n+\frac{1}{3}} \\ - 6615f_{n+\frac{2}{3}} - 756f_{n+\frac{1}{6}} - 6156f_{n+\frac{2}{6}}) \\ + h\delta_{n+1} - y_n - 3y_{n+1} + 4y_{n+\frac{1}{2}} \end{cases} = 0,(31) \\ \end{cases} \\\begin{cases} \frac{h^3}{725760} \left(-1481f_n + 457f_{n+1} + 42496f_{n+\frac{1}{2}} \\ + 120867f_{n+\frac{1}{3}} + 35805f_{n+\frac{2}{3}} + 39876f_{n+\frac{1}{6}} \\ + 3900f_{n+\frac{2}{6}} \right) + h^2\omega_{n+\frac{1}{6}} - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \end{cases} = 0,(32) \\\\ \end{cases} \\\begin{cases} \frac{h^3}{725760} \left(245f_n - 85f_{n+1} + 75008f_{n+\frac{1}{2}} \\ + 26889f_{n+\frac{1}{3}} + 21207f_{n+\frac{2}{3}} - 10380f_{n+\frac{1}{6}} \\ + 26889f_{n+\frac{1}{3}} + 21207f_{n+\frac{2}{3}} - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \end{cases} = 0,(33) \\\\ \end{cases} \\\begin{cases} \frac{h^3}{725760} \left(-11f_n + 11f_{n+1} - 1287f_{n+\frac{1}{3}} + 1287f_{n+\frac{2}{3}} \\ - 180f_{n+\frac{1}{6}} + 180f_{n+\frac{5}{6}} \right) + h^2\omega_{n+\frac{1}{2}} - 4\left(y_n + y_{n+1} \\ - 2y_{n+\frac{1}{2}} \right) \end{cases} \end{cases} = 0,(34) \\\\ \end{cases} \\\begin{cases} \frac{h^3}{725760} \left(85f_n - 245f_{n+1} - 75008f_{n+\frac{1}{2}} \\ - 21207f_{n+\frac{5}{3}} - 26889f_{n+\frac{2}{3}} - 8076f_{n+\frac{1}{6}} \\ + 10380f_{n+\frac{5}{6}} \right) + h^2\omega_{n+\frac{2}{3}} - 4\left(y_n + y_{n+1} \\ - 2y_{n+\frac{1}{2}} \right) \end{cases} = 0,(35) \\\\ \end{cases} \\\begin{cases} \frac{h^3}{725760} \left(-457f_n + 1481f_{n+1} - 42496f_{n+\frac{1}{2}} \\ - 35805f_{n+\frac{5}{3}} - 120867f_{n+\frac{2}{3}} - 3900f_{n+\frac{1}{6}} \\ - 39876f_{n+\frac{5}{6}} \right) + h^2\omega_{n+\frac{5}{3}} - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}} \right) \end{cases} \end{cases} \end{cases} = 0,(36) \\\\ \end{cases} \\\end{cases} \\\begin{cases} \frac{h^3}{26880} \left(47f_n - 1359f_{n+1} - 4352f_{n+\frac{1}{2}} + 171f_{n+\frac{1}{3}} \\ - 1035f_{n+\frac{2}{3}} - 612f_{n+\frac{1}{6}} - 6300f_{n+\frac{5}{6}} \right) + h^2\omega_{n+1} \\ - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}} \right) \end{cases} \end{cases}$$

. 2

For the one-step block hybrid method with m = 7 equally spaced off-grid points, we evaluate Equation (2) at

$$x = \left\{ x_{n+\frac{1}{8}}, \ x_{n+\frac{1}{4}}, \ x_{n+\frac{3}{8}}, \ x_{n+\frac{1}{2}}, \ x_{n+\frac{5}{8}}, \ x_{n+\frac{3}{4}}, \ x_{n+\frac{7}{8}}, \ x_{n+1} \right\}.$$

We generate the six main methods (38), (39), (40), (41), (42)

and (43) given below as:

$$\begin{cases} \frac{h^3}{88473600} \left(937f_n + 142f_{n+1} + 131885f_{n+\frac{1}{2}} + 144631f_{n+\frac{1}{4}} + 33061f_{n+\frac{3}{4}} + 53474f_{n+\frac{1}{8}} + 161518f_{n+\frac{1}{4}} + 33061f_{n+\frac{3}{4}} + 53474f_{n+\frac{1}{8}} + 161518f_{n+\frac{1}{8}} + 71398f_{n+\frac{5}{4}} + 7754f_{n+\frac{7}{8}} + 161518f_{n+\frac{1}{8}} + 71398f_{n+\frac{5}{4}} + 7754f_{n+\frac{7}{8}} + 122f_{n+\frac{1}{4}} + 14y_{n+\frac{1}{2}} - 32y_{n+\frac{1}{8}} + 1122f_{n+\frac{1}{4}} + 4742f_{n+\frac{3}{4}} + 3428f_{n+\frac{1}{8}} + 22676f_{n+\frac{3}{8}} + 10636f_{n+\frac{5}{8}} + 1148f_{n+\frac{7}{8}} + 22676f_{n+\frac{3}{8}} + 10636f_{n+\frac{5}{8}} + 1148f_{n+\frac{7}{8}} + 12863040 \left(296f_n + 199f_{n+1} + 183445f_{n+\frac{1}{2}} + 75337f_{n+\frac{1}{4}} + 46259f_{n+\frac{3}{4}} + 10858f_{n+\frac{1}{8}} + 169886f_{n+\frac{3}{8}} + 100054f_{n+\frac{3}{8}} + 10858f_{n+\frac{1}{8}} + 169886f_{n+\frac{3}{8}} + 100054f_{n+\frac{3}{8}} + 10858f_{n+\frac{1}{8}} + 169886f_{n+\frac{3}{8}} + 100054f_{n+\frac{3}{8}} - 100054f_{n+\frac{3}{8}} - 169886f_{n+\frac{3}{8}} - 18466f_{n+\frac{7}{8}} + 18466f_{n+\frac{7}{8}} + 100054f_{n+\frac{3}{8}} - 18466f_{n+\frac{7}{8}} + 128866f_{n+\frac{3}{8}} - 169886f_{n+\frac{3}{8}} - 18466f_{n+\frac{7}{8}} + 128866f_{n+\frac{3}{8}} - 169886f_{n+\frac{3}{8}} - 18466f_{n+\frac{7}{8}} + 128666f_{n+\frac{7}{8}} + 128666f_{n$$

The starters Equations (44) and (45) are obtained from Equation

(7) and are given as:

$$\begin{cases} \frac{h^{3}}{4989600} \left(-4519f_{n}+79f_{n+1}-55520f_{n+\frac{1}{2}}\right) \\ -80024f_{n+\frac{1}{4}}-13672f_{n+\frac{3}{4}}-71648f_{n+\frac{1}{8}} \\ -126304f_{n+\frac{3}{8}}-58016f_{n+\frac{5}{8}}-6176f_{n+\frac{7}{8}}\right) \\ +h\delta_{n}+3y_{n}+y_{n+1}-4y_{n+\frac{1}{2}} \end{cases} = 0, \quad (44)$$

$$\begin{cases} \frac{h^{3}}{226800} \left(8121f_{n}-209f_{n+1}-18160f_{n+\frac{1}{2}} \\ -2552f_{n+\frac{1}{4}}-4872f_{n+\frac{3}{4}}+44192f_{n+\frac{1}{8}} \\ +55584f_{n+\frac{3}{8}}+28384f_{n+\frac{5}{8}}+2912f_{n+\frac{7}{8}}\right) +h^{2}\omega_{n} \\ -4\left(y_{n}+y_{n+1}-2y_{n+\frac{1}{2}}\right) \end{cases} = 0.(45)$$

It is worth mentioning that the derivatives are generated by $\delta(x_{n+\tau_m}) = \delta_{n+\tau_m}$ and $\omega(x_{n+\tau_m}) = \omega_{n+\tau_m}$

$$\begin{cases} \frac{h^{3}}{5109350400} \left(371725f_{n} - 65591f_{n+1} - 43219630f_{n+\frac{1}{2}} - 31649144f_{n+\frac{1}{4}} - 10970288f_{n+\frac{3}{4}} + 5710568f_{n+\frac{1}{8}} - 44180864f_{n+\frac{3}{4}} - 20167160f_{n+\frac{5}{8}} - 2191216f_{n+\frac{7}{8}} \right) + h\delta_{n+\frac{1}{8}} + \frac{1}{2} \left(5y_{n} + y_{n+1} - 6y_{n+\frac{1}{2}} \right) \\ \left(\frac{h^{3}}{79833600} \left(2221f_{n} + 65f_{n+1} + 10930f_{n+\frac{1}{2}} + 515828f_{n+\frac{1}{4}} + 2524f_{n+\frac{3}{4}} + 156736f_{n+\frac{1}{8}} + 149920f_{n+\frac{3}{8}} - 6016f_{n+\frac{5}{8}} - 608f_{n+\frac{7}{8}} \right) \\ + 149920f_{n+\frac{3}{8}} - 6016f_{n+\frac{5}{8}} - 608f_{n+\frac{7}{8}} \right) \\ + h\delta_{n+\frac{1}{4}} + 2 \left(y_{n} - y_{n+\frac{1}{2}} \right) \end{cases} = 0, \quad (47)$$

$$\left(\frac{h^{3}}{5109350400} \left(128801f_{n} + 39349f_{n+1} + 40485970f_{n+\frac{1}{2}} + 30702112f_{n+\frac{1}{4}} + 9998264f_{n+\frac{3}{4}} + 7176736f_{n+\frac{1}{8}} + 59573336f_{n+\frac{3}{8}} + 22437424f_{n+\frac{5}{8}} + 2430808f_{n+\frac{7}{8}} \right) + h\delta_{n+\frac{3}{8}} + \frac{1}{2} \left(3y_{n} - y_{n+1} \right) \\ - 2y_{n+\frac{1}{2}} \right) \end{cases} = 0, \quad (48)$$

$$\left(\frac{h^{3}}{4989600} \left(79f_{n} + 79f_{n+1} + 70430f_{n+\frac{1}{2}} + 19504f_{n+\frac{3}{4}} + 19504f_{n+\frac{3}{4}} + 44736f_{n+\frac{1}{8}} + 44416f_{n+\frac{3}{8}} + 4736f_{n+\frac{7}{8}} \right) + h\delta_{n+\frac{1}{2}} + y_{n} - y_{n+1} \right) \right\} = 0, \quad (49)$$

$$\begin{cases} \frac{h^3}{5109350400} \left(39349f_n + 128801f_{n+1} + 40485970f_{n+\frac{1}{2}} + 9998264f_{n+\frac{1}{4}} + 30702112f_{n+\frac{1}{2}} + 2430808f_{n+\frac{1}{4}} + 22437424f_{n+\frac{1}{2}} + 59573336f_{n+\frac{5}{2}} + 22437424f_{n+\frac{1}{2}} + 59573336f_{n+\frac{5}{2}} + 22437424f_{n+\frac{1}{2}} + 59573336f_{n+\frac{5}{2}} + 2524f_{n+\frac{1}{2}} + 515828f_{n+\frac{1}{2}} - 608f_{n+\frac{1}{4}} + 2524f_{n+\frac{1}{4}} + 515828f_{n+\frac{1}{2}} - 608f_{n+\frac{1}{4}} + 2524f_{n+\frac{1}{4}} + 515828f_{n+\frac{1}{2}} - 608f_{n+\frac{1}{4}} + 2(y_{n+\frac{1}{2}} - y_{n+1}) \\ \\ \frac{h^3}{5109350400} \left(-65591f_n + 371725f_{n+1} - 43219630f_{n+\frac{1}{2}} - 10970288f_{n+\frac{1}{4}} - 31649144f_{n+\frac{2}{2}} - 2191216f_{n+\frac{1}{4}} - 31649144f_{n+\frac{2}{3}} - 2191216f_{n+\frac{1}{4}} - 31649144f_{n+\frac{2}{3}} - 2191216f_{n+\frac{1}{4}} - 31649144f_{n+\frac{2}{3}} - 126304f_{n+\frac{2}{5}} + \frac{1}{2}\left(-y_n - 5y_{n+1} \right) \\ \\ = 0, \quad (52) \\ \\ \frac{h^3}{4989600} \left(79f_n - 4519f_{n+1} - 55520f_{n+\frac{1}{2}} - 13672f_{n+\frac{1}{4}} - 80024f_{n+\frac{2}{3}} - 71648f_{n+\frac{2}{3}} \right) \\ \\ = 0, \quad (53) \\ \\ \frac{h^3}{29030400} \left(-30529f_n + 7201f_{n+1} + 2708640f_{n+\frac{2}{3}} + 4277938f_{n+\frac{1}{4}} + 667598f_{n+\frac{3}{4}} \right) \\ \\ + 1189482f_{n+\frac{1}{4}} + 1519394f_{n+\frac{3}{4}} + 486814f_{n+\frac{5}{4}} \\ + 59862f_{n+\frac{1}{4}} \right) + h^2\omega_{n+\frac{1}{4}} - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \\ \\ \\ \frac{h^3}{907200} \left(107f_n - 3f_{n+1} + 43480f_{n+\frac{1}{4}} \\ + 32286f_{m+\frac{1}{4}} + 11666f_{n+\frac{2}{4}} - 5816f_{n+\frac{1}{4}} \\ + 102248f_{n+\frac{3}{4}} + 38088f_{n+\frac{5}{4}} \\ + 4024f_{n+\frac{7}{4}} \right) + h^2\omega_{n+\frac{1}{4}} - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{4}}\right) \\ \\ \\ \end{array} \right\} = 0, (55) \\ \\ \end{array}$$

$$\begin{cases} \frac{h^{3}}{29030400} \left(-3873f_{n} + 3137f_{n+1} + 2224480f_{n+\frac{1}{2}} \right) \\ -605134f_{n+\frac{1}{4}} + 515406f_{n+\frac{3}{4}} - 86486f_{n+\frac{1}{8}} \\ +640098f_{n+\frac{3}{8}} + 843998f_{n+\frac{5}{8}} + 97174f_{n+\frac{7}{8}} \\ +h^{2}\omega_{n+\frac{3}{8}} - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \end{cases} = 0,(56)$$

$$\frac{45360}{45360} \left(-f_n + f_{n+1} - 608f_{n+\frac{1}{4}} + 608f_{n+\frac{3}{4}} - 192f_{n+\frac{1}{8}} - 1984f_{n+\frac{3}{8}} + 1984f_{n+\frac{5}{8}} + 192f_{n+\frac{7}{8}} \right) = 0,(57)$$

$$+ h^2 \omega_{n+\frac{1}{2}} - 4 \left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}} \right)$$

$$\begin{cases} \frac{h^3}{29030400} \left(-3137f_n + 3873f_{n+1} - 2224480f_{n+\frac{1}{2}} \right) \\ -515406f_{n+\frac{1}{4}} + 605134f_{n+\frac{3}{4}} - 97174f_{n+\frac{1}{8}} \\ -843998f_{n+\frac{3}{8}} - 640098f_{n+\frac{5}{8}} + 86486f_{n+\frac{7}{8}} \right) \\ + h^2\omega_{n+\frac{5}{8}} - 4\left(y_n + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \end{cases} = 0,(58)$$

$$\begin{cases} \frac{h^{3}}{907200} \left(3f_{n} - 107f_{n+1} - 43480f_{n+\frac{1}{2}} - 11666f_{n+\frac{1}{4}}\right) \\ - 32286f_{n+\frac{3}{4}} - 4024f_{n+\frac{1}{8}} - 38808f_{n+\frac{3}{8}} \\ - 102248f_{n+\frac{5}{8}} + 5816f_{n+\frac{7}{8}}\right) + h^{2}\omega_{n+\frac{3}{4}} \\ - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{2}}\right) \end{cases} =$$

$$\left. \frac{h^{3}}{29030400} \left(-7201f_{n} + 30529f_{n+1} - 2708640f_{n+\frac{1}{2}} - 667598f_{n+\frac{1}{4}} - 4277938f_{n+\frac{3}{4}} - 59862f_{n+\frac{1}{8}} - 486814f_{n+\frac{3}{8}} - 1519394f_{n+\frac{5}{8}} - 1189482f_{n+\frac{7}{8}} \right) + h^{2}\omega_{n+\frac{7}{8}} - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{3}}\right) + h^{2}\omega_{n+\frac{7}{8}} - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{3}}\right)$$

$$\frac{h^{3}}{226800} \left(209f_{n} - 8121f_{n+1} + 18160f_{n+\frac{1}{2}} + 4872f_{n+\frac{1}{4}} + 2552f_{n+\frac{3}{4}} - 2912f_{n+\frac{1}{8}} - 28384f_{n+\frac{3}{8}} - 55584f_{n+\frac{5}{8}} - 44192f_{n+\frac{7}{8}} \right) \\
+ h^{2}\omega_{n+1} - 4\left(y_{n} + y_{n+1} - 2y_{n+\frac{1}{2}}\right) + (61)$$

For brevity, we have omitted the schemes for m > 7.

2.3. Analysis of the methods

We report the findings from our analysis of the characteristics of our proposed method with evenly spaced off-grid collocation nodes in this section. We focus in particular on the zero stability analysis and truncation error.

2.3.1. Local Truncation Errors and Order

We express our main methods and starters for third order IVPs in terms of a linear operator \mathcal{L} defined as

$$\sum_{j=0}^{1} \alpha_{j} y(x_{n} + jh) + \alpha_{p_{\frac{1}{2}}} y(x_{n} + p_{\frac{1}{2}}h) \\ - h^{3} (\sum_{i=0}^{1} \beta_{i} y^{'''}(x_{n} + ih) + \sum_{\nu=1}^{m} \beta_{p_{\nu}} y^{'''}(x_{n} + p_{\nu}h)), \\ \sum_{j=0}^{1} \alpha_{j} y(x_{n} + jh) - hy'(x_{n}) \\ - h^{3} (\sum_{i=0}^{1} \beta_{i} y^{'''}(x_{n} + ih) + \sum_{\nu=1}^{m} \beta_{p_{\nu}} y^{'''}(x_{n} + p_{\nu}h)), \\ \sum_{j=0}^{1} \alpha_{j} y(x_{n} + jh) - hy''(x_{n}) \\ - h^{3} (\sum_{i=0}^{1} \beta_{i} y^{'''}(x_{n} + ih) + \sum_{\nu=1}^{m} \beta_{p_{\nu}} y^{'''}(x_{n} + p_{\nu}h)), \end{cases}$$

where $\zeta = 1, 2, ..., m$.

Assuming that $y(x_n)$ is sufficiently differentiable, we can expand the terms $y(x_n + jh)$, $y(x_n + p_{\zeta}h)$, $y'''(x_n + p_{\nu}h)$ and $y'''(x_n + jh)$ as a Taylor series about the point x_n to obtain the expression

$$L[y(x);h] = \hat{C}_0 y(x_n) + \hat{C}_1 h y'(x_n) + \dots + \hat{C}_p h^p y^p(x_n) + \dots (63)$$

= 0,(59)where $\hat{C}_0, \hat{C}_1, \dots, \hat{C}_p$ are constant vectors. As stated in Henrici [1], we say that the method has order *p* if

$$\hat{C}_0, \hat{C}_1, \cdots, \hat{C}_p, \hat{C}_{p+1}, \text{ and } \hat{C}_{p+2} = 0 \quad C_{p+3} \neq 0.$$

The vector \hat{C}_{p+3} is called the error constant and $\hat{C}_{p+3}h^{p+3}y^{(p+3)}(x_n)$ the principal part of the local truncation error at the end point x_n . It is standard from our computations that our methods have order p > 1 and with comparatively small error constants.

Consequently, for our method with m = 3, we have

$$\frac{h^{8}y^{(8)}(x)}{125829120} + O(h^{9}), \\
\frac{h^{8}y^{(8)}(x)}{125829120} + O(h^{9}), \\
\frac{h^{8}y^{(8)}(x)}{5160960} + O(h^{9}), \\
-\frac{h^{8}y^{(8)}(x)}{33030144} + O(h^{9})$$
(64)

For our method with m = 5 and m = 7, we have

$$\frac{13h^{10}y^{(10)}(x)}{10970982973440} + O(h^{11}), \\ - \frac{h^{10}y^{(10)}(x)}{548549148672} + O(h^{11}), \\ - \frac{h^{10}y^{(10)}(x)}{548549148672} + O(h^{11}), \\ \frac{13h^{10}y^{(10)}(x)}{10970982973440} + O(h^{11}), \\ \frac{h^{10}y^{(10)}(x)}{4180377600} + O(h^{11}), \\ - \frac{h^{10}y^{(10)}(x)}{179159040} + O(h^{11}) \\ \end{array} \right] = \mathcal{L}[Y(x_n); h], \quad (65)$$

$$\frac{7h^{12}y^{(12)}(x)}{2597596220620800} + O(h^{13}), \\ - \frac{h^{12}y^{(12)}(x)}{25955487255461888} + O(h^{13}), \\ - \frac{h^{12}y^{(12)}(x)}{2597596220620800} + O(h^{13}), \\ - \frac{h^{12}y^{(12)}(x)}{50226958172160} + O(h^{13}) \\ - \frac{h^{12}y^{(12)}(x)}{50226958172160} + O(h^{13}) \\ - \frac{h^{12}y^{(12)}(x)}{50226958172160} + \frac{h^{12}y^{($$

2.3.2. Zero Stability

The starters and primary methods are represented in matrixvector form as

$$A_m^{(1)}Y_{\lambda+1} = A_m^{(0)}Y_{\lambda} + h^3 [B_m^{(1)}F_{\lambda+1} + B_m^{(0)}F_{\lambda}] + h^2 C_m^{(0)}\omega_{\lambda} + h D_m^{(0)}\delta_{\lambda}, (67)$$

T

where

$$Y_{\lambda+1} = (\cdots, y_{n+\frac{1}{2}}, \cdots, y_{n+1})^T,$$

$$Y_{\lambda} = (\cdots, y_{n-\frac{1}{2}}, \cdots, y_n)^T,$$

$$F_{\lambda+1} = (\cdots, f_{n+y_{n+\frac{1}{2}}}, \cdots, f_{n+2})^T,$$

$$F_{\lambda} = (\cdots, f_{n+y_{n-\frac{1}{2}}}, \cdots, f_n)^T,$$

$$\delta_{\lambda} = (\cdots, \delta_{n-\frac{1}{2}}, \cdots, \delta_n)^T,$$

$$\omega_{\lambda} = (\cdots, \omega_{n-\frac{1}{2}}, \cdots, \omega_n)^T,$$

 $A_m^{(1)}, A_m^{(0)}, B_m^{(1)}, B_m^{(0)}, C_m^{(0)}$ and $D_m^{(0)}$ are square matrices.

For m = 3 we obtain we obtain the square matrices

$$A_{3}^{(1)} = \begin{pmatrix} 1 & -\frac{3}{4} & 0 & \frac{1}{8} \\ 0 & -\frac{3}{4} & 1 & -\frac{3}{8} \\ 0 & 4 & 0 & -1 \\ 0 & -8 & 0 & 4 \end{pmatrix}, A_{0}^{(0)} = \begin{pmatrix} 0 & 0 & 0 & -\frac{3}{8} \\ 0 & 0 & 0 & \frac{1}{8} \\ 0 & 0 & 0 & -3 \\ 0 & 0 & 0 & 4 \end{pmatrix},$$
$$B_{3}^{(1)} = \begin{pmatrix} -\frac{43}{15360} & \frac{21}{5120} & -\frac{13}{15360} & \frac{1}{30720} \\ -\frac{13}{15360} & -\frac{21}{5120} & -\frac{13}{15360} & -\frac{1}{30720} \\ -\frac{53}{1260} & -\frac{13}{55} & -\frac{11}{1260} & \frac{1}{10080} \\ \frac{3}{10} & \frac{1}{15} & \frac{1}{18} & -\frac{1}{240} \end{pmatrix},$$
$$B_{3}^{(0)} = \begin{pmatrix} 0 & 0 & 0 & \frac{1}{30720} \\ 0 & 0 & 0 & -\frac{30720}{10080} \\ 0 & 0 & 0 & -\frac{41}{10080} \\ 0 & 0 & 0 & \frac{59}{720} \end{pmatrix}, C_{3}^{(0)} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$
and

For m = 5 we obtain

	(1	0	$-\frac{5}{9}$	0	0	$\frac{1}{9}$)
	0	1	$-\frac{8}{9}$	0	0	$\frac{1}{9}$	
A (1) _	0	0	$-\frac{8}{9}$	1	0	$-\frac{2}{9}$	
$A_{5} =$	0	0	$-\frac{5}{9}$	0	1	$-\frac{5}{9}$	
	0	0	4	0	0	-1	
	0	0	-8	0	0	4	J

 $B_5^{(1)} =$

7891	3835	4405	689	1591	97
6531840	1306368	1952552	653184	6531840	19595520
/9/	13513	1/9	6583	419	13
1632960	6531840	76545	6531840	1632960	3919104
$-\frac{11}{1632960}$	$-\frac{0000}{6531840}$	$-\frac{17}{76545}$	$-\frac{13313}{6531840}$	$-\frac{757}{1632960}$	$-\frac{191}{19595520}$
1591	689	4405	3835	7891	103
6531840	653184	1959552	1306368	6531840	4898880
$-\frac{513}{22400}$	$-\frac{63}{2560}$	$-\frac{1}{42}$	$-\frac{27}{2504}$		11
22400	2560	17^{42}	3284	3200	268800
$\frac{15}{64}$	$\frac{0}{1792}$	$\frac{17}{105}$	$-\frac{37}{8960}$	$\frac{31}{2240}$	$-\frac{1}{26880}$

For brevity, we omit the square matrices for the method m = 7. The matrix system (67) is reduced to

$$A_m^{(1)}Y_{\lambda+1} = A_m^{(0)}Y_{\lambda}.$$
(68)

The zero stability is established from the nature of the zeros of the first characteristic polynomial defined as

$$\rho(R) = det[A^{(1)} - A^{(0)}]. \tag{69}$$

A method is said to be zero stable if the roots of $\rho(R)$ satisfy $|R_j| \le 1$ and all the roots with $|R_j| = 1$ have multiplicity that does not exceed 2. For our method with m = 3, we obtain the first characteristic polynomial given as

$$\rho(R) = -8R^3(1+R),\tag{70}$$

which has principal root $| R_0 |= 1$ and spurious roots $R_j = 0$, j = 1(1)3. For our method with m = 5, we obtain the first characteristic polynomial given as

$$\rho(R) = 8R^{3}(1+R), \tag{71}$$

which has principal root $| R_0 |= 1$ and spurious roots $R_j = 0$, j = 1(1)5. For our method with m = 7, we obtain the first characteristic polynomial given as

$$\rho(R) = -8R'(1+R), \tag{72}$$

which has principal root $|R_0| = 1$ and spurious roots $R_j = 0$, j = 1(1)7. In general, the first characteristic polynomial of our class of hybrid block methods alternates between

$$\rho(R) = -R^m(R+1),$$
(73)

and

$$\rho(R) = R^m (R+1).$$
(74)

Following Jator [17], our class of hybrid block methods are zero stable, consistent and convergent.

3. Numerical Problems & Discussions

In this section, we have tested the performance of our methods on both scalar and system ODEs of the linear and nonlinear types. All results in this study were compared with the theoretical solutions, hence, the Absolute Errors (AE) and Maximum Absolute Errors (MAE) obtained. We compare our methods with the Generalized Linear Block Method (GLBM) of order 5 in Ref. [3], the Direct Two-Point Four-Step Variable Step (D2P4VS) methods of order 6 and 7 in Ref. [4], the Implicit Three-Point Block Method of order 9 (ITPBMO9) in Ref. [5], and the Boundary Value Methods (BVM5) of order 6 in Ref. [11]. All computations were implemented using Mathematica 13.

Ta	Table 1. Comparison of the MAE of $m = 3$ and BVM5 (Problem 1)		
N	<i>m</i> = 3	BVM5	
10	2.91×10^{-9}	5.07×10^{-6}	
20	4.65×10^{-11}	1.16×10^{-7}	
40	7.39×10^{-13}	2.07×10^{-9}	
80	1.84×10^{-14}	3.35×10^{-11}	
160	1.65×10^{-13}	5.92×10^{-13}	
320	2.52×10^{-12}	8.28×10^{-14}	

3.1. Problem 1

Consider the non-linear third order IVP

$$y^{'''}(x) + 2e^{-3y(x)} = 4(1+x)^{-3},$$
 (75)

which was solved in Ref. [11], with the initial conditions, interval of integration and exact solution given as

$$y(0) = 0, y'(0) = 1, y''(0) = -1,$$

 $x \in [0, 1],$
 $Exact : y(x) = \ln(1 + x).$

In this example, a comparison is made between our method m = 3 and the BVM5 in Ref. [11], both orders 5 and 6 respectively. In Table 1, we compare the MAE of our method m = 3 with the BVM5 at different values of N. It is evident from the results that our method m = 3 of order 5 performs favourably well against the BVM5 of order 6. Our method m = 3 outperforms the BVM5 for all the values of N provided except at 320.

Table 2 display the MAE for our method m = 3, 5 and 7. The results were computed using different values for N.

3.2. Problem 2

Consider the singularly perturbed third order BVP

$$-\epsilon y^{'''}(x) + 4y'(x) = 1, \ x \in [0, 1], \tag{76}$$

which was solved in Refs. [6, 11], with the boundary conditions, domain of integration and exact solution given as

$$y(0) = y'(0) = y'(1) = 1,$$

$$Exact: y(x) = \frac{x}{4} - \frac{c_1\sqrt{\epsilon}}{2}\exp\left(\frac{-2x}{\sqrt{\epsilon}}\right) + \frac{c_2\sqrt{\epsilon}}{2}\exp\left(\frac{2x}{\sqrt{\epsilon}}\right) + c_3,$$

where

$$c_1 = \frac{3}{4} - \frac{3}{4\left(1 + \exp\left(\frac{2}{\sqrt{\epsilon}}\right)\right)}$$

Т	able 2. MAE methods m	i = 3, 5 and 7 with vari	ous N (Problem 1)
N	<i>m</i> = 3	<i>m</i> = 5	m = 7
6	5.93×10^{-8}	1.52×10^{-10}	1.69×10^{-6}
10	2.91×10^{-9}	2.76×10^{-12}	2.28×10^{-8}
12	9.84×10^{-10}	6.52×10^{-13}	4.45×10^{-9}
20	4.65×10^{-11}	7.97×10^{-15}	3.73×10^{-11}
24	1.56×10^{-11}	5.53×10^{-15}	6.49×10^{-12}
40	7.39×10^{-13}	5.76×10^{-15}	1.08×10^{-14}
48	2.32×10^{-13}	1.34×10^{-14}	1.71×10^{-14}
80	1.84×10^{-14}	2.41×10^{-14}	8.40×10^{-14}
96	1.76×10^{-14}	7.64×10^{-14}	1.84×10^{-13}
160	1.65×10^{-13}	1.36×10^{-13}	5.79×10^{-14}
192	5.48×10^{-14}	1.58×10^{-13}	1.60×10^{-13}
320	2.52×10^{-12}	3.19×10^{-12}	1.01×10^{-12}

$$c_{2} = \frac{3}{4\left(1 + \exp\left(\frac{2}{\sqrt{\epsilon}}\right), \frac{1}{\sqrt{\epsilon}}\right)},$$

$$c_{3} = \frac{8 + 8\exp\left(\frac{2}{\sqrt{\epsilon}}\right) - 3\sqrt{\epsilon} + 3\sqrt{\epsilon}\exp\left(\frac{2}{\sqrt{\epsilon}}\right)}{8\left(1 + \exp\left(\frac{2}{\sqrt{\epsilon}}\right)\right)}.$$

The singularly perturbed third-order BVP has been solved with a variety of values ϵ , and it is shown that the results are still acceptable as $\epsilon \to 0$. In the Tables 3 and 4, the MAE of our method m = 3 of order 5 is compared with the BVM5 of order 6. From the results presented, our method m = 3 outperforms the BVM5 for $\epsilon = 0.01$. Although, our methods perform better with a decreased number of steps, for example, our methods m = 3, 5 and 7 outperform those presented in Ref. [11] for smaller values of N, the BVM5 outperforms our methods at N = 320 and with $\epsilon = 0.1$.

For our methods, m = 3, 5 and 7, the numerical solutions were contrasted with the theoretical and the MAE presented in Table 5 and 6.

3.3. Problem 3

Consider the third order IVP

$$y^{'''} + 2y^{''} - 9y^{'} - 18y = -18x^2 - 18x + 22,$$
 (77)

with the initial conditions, domain of integration and exact solution given as

Table 3.	able 3. Comparison of the MAE $m = 3$ and BVM5 (Problem 2, $\epsilon = 0.0$	
N	m = 3	BVM5
10	3.7×10^{-6}	2.6×10^{-3}
20	8.4×10^{-8}	1.1×10^{-4}
40	1.5×10^{-9}	6.9×10^{-6}
80	2.3×10^{-11}	4.3×10^{-7}
160	3.7×10^{-13}	2.7×10^{-12}
320	1.3×10^{-14}	1.7×10^{-9}

Table 4. Comparison of the MAE m = 3 and BVM5 (Problem 2, $\epsilon = 0.1$)

BVM5	m = 3	N
2.4×10^{-5}	$1.8 imes 10^{-8}$	10
4.8×10^{-7}	2.9×10^{-10}	20
8.0×10^{-9}	4.6×10^{-12}	40
1.3×10^{-10}	7.7×10^{-14}	80
2.0×10^{-12}	1.4×10^{-14}	160
1.8×10^{-14}	1.8×10^{-13}	320

$$y(0) = -2, y'(0) = -8, y''(0) = 12, x \in [0, b]$$

Exact :
$$y(x) = -2e^{-3x} + e^{-2x} + x^2 - 1$$
.

Tables 7 and 8 display the AE at the endpoints x = 1 and x = 4 of the integration domain. The results were computed using different values for *N*.

We compare our methods m = 3 of order 5, with the two-point four-step implicit block method (D2P4VS) of orders 6 and 7 presented in Ref. [4].

In the Tables 9 and 10, we provide the MAE within the integration domain $x \in [0, b]$. From the results displayed in Table 9, our methods outperform the D2P4VS in a lesser number of steps with b = 1. In the Table 10, our method m = 3 outperforms the D2P4VS.

3.4. Problem 4

Consider the BVP with mixed boundary conditions

$$y'''(x) - \frac{1}{\sqrt{1+x}} + 2y(x) = f(x), \tag{78}$$

Table :	5. MAE method $m = 3$, 5 and 7 with various	N (Problem 2, $\epsilon = 0.01$)	Table 6	6. MAE methods $m =$	3, 5 and 7 with various	N (problem 2, $\epsilon = 0.1$)
Ν	<i>m</i> = 3	<i>m</i> = 5	<i>m</i> = 7	Ν	<i>m</i> = 3	<i>m</i> = 5	<i>m</i> = 7
6	3.9×10^{-5}	8.01×10^{-7}	1.10×10^{-8}	6	3.5×10^{-7}	6.74×10^{-10}	8.93×10^{-13}
10	3.7×10^{-6}	2.63×10^{-8}	1.27×10^{-10}	10	1.8×10^{-8}	1.25×10^{-11}	6.22×10^{-15}
12	1.4×10^{-6}	7.01×10^{-9}	2.33×10^{-11}	12	6.1×10^{-9}	2.95×10^{-12}	3.11×10^{-15}
20	8.4×10^{-8}	1.46×10^{-10}	1.75×10^{-13}	20	2.9×10^{-10}	4.51×10^{-14}	4.22×10^{-15}
24	2.9×10^{-8}	3.54×10^{-11}	3.22×10^{-14}	24	9.9×10^{-11}	1.31×10^{-14}	4.44×10^{-15}
40	1.5×10^{-9}	6.31×10^{-13}	8.88×10^{-16}	40	4.6×10^{-12}	1.58×10^{-14}	7.33×10^{-15}
48	4.9×10^{-10}	1.51×10^{-13}	2.58×10^{-14}	48	1.5×10^{-12}	3.33×10^{-15}	7.33×10^{-15}
80	2.3×10^{-11}	3.33×10^{-15}	4.00×10^{-15}	80	7.7×10^{-14}	2.89×10^{-14}	2.02×10^{-14}
96	7.9×10^{-12}	7.55×10^{-15}	1.55×10^{-15}	96	4.8×10^{-14}	4.06×10^{-14}	1.26×10^{-12}
160	3.7×10^{-13}	7.59×10^{-14}	6.88×10^{-15}	160	1.4×10^{-14}	2.91×10^{-14}	5.13×10^{-14}
192	1.2×10^{-13}	9.99×10^{-15}	7.37×10^{-13}	192	1.1×10^{-13}	7.86×10^{-13}	4.57×10^{-14}
320	1.3×10^{-14}	8.00×10^{-15}	5.45×10^{-12}	320	1.8×10^{-13}	2.25×10^{-12}	5.20×10^{-14}

which was solved in Ref. [11], with the mixed boundary conditions, domain of integration and exact solution given as

$$y(0) = y'(0) = 1, y(1) - y'(1) = 0, x \in [0, 1],$$

 $Exact: y(x) = \frac{1}{2}x^{3}.$

In this problem, f(x) is deduced from the exact solution. All methods solve this problem accurately and are consequently exact to machine precision. Table 11 shows the computational comparison of the different methods for this problem.

In Table 12, we display the MAE within the integration domain of the problem for m = 3, 5 and 7. The results were computed using different values of N.

3.5. Problem 5

Consider the third order IVP

$$y^{'''} + e^x = 0, (79)$$

which was solved in Ref. [3], with the initial conditions, domain of integration and exact solution given as

$$y(0) = 1, y'(0) = -1, y''(0) = 3, x \in [0, 1],$$

 $Exact : y(x) = 2x^2 - e^x + 2.$

Tables 13 and 14 show the AE at selected points in the integration domain for the problem solved with step size $h = \frac{1}{10}$

160	1.4×10^{-14}	2.91×10^{-14}	5.13×10^{-14}
192	1.1×10^{-13}	7.86×10^{-13}	4.57×10^{-14}
320	1.8×10^{-13}	2.25×10^{-12}	5.20×10^{-14}
	Table 7 AE at the	e end of interval (Problem	(3, r-1)
	Tuble 7. THE ut un		
Ν	m = 3	m = 5	$\frac{m}{m} = 7$
N 8	m = 3 2.456667 × 10 ⁻⁷	$\frac{1}{m} = 5$ 6.023624 × 10 ⁻¹¹	$\frac{m = 7}{1.010961 \times 10^{-14}}$
<u>N</u> 8 10	m = 3 2.456667 × 10 ⁻⁷ 6.421349 × 10 ⁻⁸	1000000000000000000000000000000000000	$\frac{m}{1.010961 \times 10^{-14}}$ 1.082454×10^{-15}

3, $x = 4$)	e end of interval (Problem	Table 8. AE at the	
<i>m</i> = 7	<i>m</i> = 5	<i>m</i> = 3	N
7.038723×10^{-13}	1.406541×10^{-8}	1.901826×10^{-4}	59
3.974098×10^{-15}	2.236027×10^{-10}	8.512268×10^{-6}	99
5.803680×10^{-16}	4.797736×10^{-11}	2.683452×10^{-6}	120

and compared with the theoretical solution. In the Table 13, we compare the AE of our methods m = 3 against the GLBM in Ref. [3]. It is instructive to note that both methods are of order 5. It is easy to see that our method m = 3 performs relatively well against the GLBM. From Table 13, both methods perform slightly better than the other at 5 selected points each within the domain of integration. The AE of our methods m = 5 and m = 7 is presented in Table 14.

Ta	Table 9. Comparison of MAE $m = 3$ and D2P4VS (problem 3, $x = 1$)		
Ν	<i>m</i> = 3	Ν	D2P4VS
8	2.46×10^{-7}	41	9.33×10^{-7}
10	6.42×10^{-8}	54	7.82×10^{-8}
22	5.64×10^{-10}	64	8.16×10^{-10}

Table	10. Comparison of MAE $m = 3$	and D2P4VS (problem 3, $x = 4$)	
N	<i>m</i> = 3	D2P4VS	
41	1.41×10^{-8}	2.26×10^{-6}	
54	6.42×10^{-8}	7.82×10^{-8}	
64	5.64×10^{-10}	8.16×10^{-10}	

Table 11. Comparison of the MAE methods m = 3 and BVM5 (Problem 4)

D V IVIJ	m = 3	11
1.38×10^{-15}	1.55×10^{-15}	10
4.44×10^{-16}	4.44×10^{-15}	20
2.87×10^{-14}	2.44×10^{-15}	40
6.86×10^{-14}	1.65×10^{-14}	80
2.45×10^{-14}	7.07×10^{-14}	160
5.23×10^{-14}	1.94×10^{-14}	320

3.6. Problem 6

Consider the nonlinear third order IVP

$$y^{'''} - xy^{''} + x^2y^2 = x\sin x - \cos x + x^2\sin^2 x,$$
 (80)

which was solved in Ref. [3], with the initial conditions, domain of integration and exact solution given as

$$y(0) = 0, y'(0) = 1, y''(0) = 0, x \in [0, 1],$$

$Exact: y(x) = \sin x.$

In this example, a comparison is made between our method m = 3 and the GLBM in Ref. [3]. In Table 15, we compare the AE of our method m = 3 with the GLBM at selected points within the integration domain of the problem solved with step size $h = \frac{1}{10}$. It is evident from the results that our method m = 3 performs relatively well against the GLBM. Our scheme m = 3 slightly outperforms the GLBM at the first five selected points in the Table 15, while the GLBM marginally outperforms our method m = 3 at the points x = 0.6, 0.7, 0.8, 0.9 and 1.0. The AE of our methods m = 5 and m = 7 is presented in Table 16.

Table 12. MAE methods $m = 3, 5$ and 7 with various N (Problem 4)				
N	m = 3	m = 5	m = 7	
6	1.10×10^{-16}	1.33×10^{-15}	8.22×10^{-18}	
10	1.55×10^{-15}	1.44×10^{-15}	1.57×10^{-16}	
12	1.22×10^{-15}	1.77×10^{-15}	1.17×10^{-15}	
20	4.44×10^{-15}	1.44×10^{-15}	4.00×10^{-15}	
24	8.10×10^{-16}	8.33×10^{-16}	5.55×10^{-16}	
40	2.44×10^{-15}	1.11×10^{-15}	3.33×10^{-15}	
48	1.37×10^{-16}	4.11×10^{-15}	8.22×10^{-15}	
80	1.65×10^{-14}	1.99×10^{-14}	9.21×10^{-15}	
96	3.15×10^{-16}	1.37×10^{-14}	9.38×10^{-15}	
160	7.07×10^{-14}	3.21×10^{-14}	6.36×10^{-14}	
192	4.15×10^{-16}	9.03×10^{-14}	2.89×10^{-14}	
320	1.94×10^{-14}	1.22×10^{-13}	2.10×10^{-13}	

Table 13. Comparison of the AE methods $m = 3$ and GLBM (Problem 5)			
x	<i>m</i> = 3	GLBM	
0.1	1.617997×10^{-14}	1.620926×10^{-14}	
0.2	6.611369×10^{-14}	$6.605827 imes 10^{-14}$	
0.3	1.528079×10^{-13}	1.527667×10^{-13}	
0.4	2.795853×10^{-13}	2.795542×10^{-13}	
0.5	4.501185×10^{-13}	4.501954×10^{-13}	
0.6	6.684659×10^{-13}	6.683543×10^{-13}	
0.7	9.391131×10^{-13}	9.392487×10^{-13}	
0.8	1.267017×10^{-12}	1.266987×10^{-12}	
0.9	1.657656×10^{-12}	1.657785×10^{-12}	
1.0	2.117086×10^{-12}	2.117417×10^{-12}	

3.7. Problem 7

Consider the linear system

Table 14. AE at selected points for method $m = 5$ and $m = 7$ (Problem 5)			Tab
x	<i>m</i> = 5	<i>m</i> = 7	x
0.1	2.667406×10^{-19}	3.083489×10^{-24}	0.1
0.2	1.091358×10^{-18}	1.262510×10^{-23}	0.2
0.3	2.525817×10^{-18}	2.924209×10^{-23}	0.3
0.4	4.627550×10^{-18}	5.361661×10^{-23}	0.4
0.5	7.460028×10^{-18}	8.650259×10^{-23}	0.5
0.6	1.109340×10^{-17}	1.287332×10^{-22}	0.6
0.7	1.560518×10^{-17}	1.812293×10^{-22}	0.7
0.8	2.108106×10^{-17}	2.450086×10^{-22}	0.8
0.9	2.761572×10^{-17}	3.211957×10^{-22}	0.9
1.0	$3.531380 imes 10^{-17}$	4.110336×10^{-22}	1.0

Table 15. Comparison of the AE methods m = 3 and GLBM (Problem 6)

x	m = 3	GLBM
0.1	6.6369×10^{-16}	6.661338×10^{-16}
0.2	3.9126×10^{-15}	3.913536×10^{-15}
0.3	1.23524×10^{-14}	1.243450×10^{-14}
0.4	2.86685×10^{-14}	2.886580×10^{-14}
0.5	5.59501×10^{-14}	$5.601075 imes 10^{-14}$
0.6	9.70193×10^{-14}	9.692247×10^{-14}
0.7	1.55024×10^{-13}	1.546541×10^{-13}
0.8	2.33561×10^{-13}	2.325917×10^{-13}
0.9	3.35928×10^{-13}	3.346212×10^{-13}
1.0	4.65962×10^{-13}	4.644063×10^{-12}

$$y_{1}^{'''}(x) = \frac{1}{68}(817y_{1} + 1393y_{2} + 448y_{3})$$
$$y_{2}^{'''}(x) = -\frac{1}{68}(1141y_{1} + 2837y_{2} + 896y_{3}) , \quad (81)$$
$$y_{3}^{'''}(x) = \frac{1}{136}(3059y_{1} + 4319y_{2} + 1592y_{3})$$

which was solved in Ref. [11], with the initial conditions

Table 1	16. AE at selected points for met	hod $m = 5$ and $m = 7$ (Problem 6)
x	<i>m</i> = 5	<i>m</i> = 7
0.1	2.44337×10^{-18}	2.44337×10^{-18}
0.2	9.39821×10^{-19}	9.39821×10^{-19}
0.3	2.65894×10^{-17}	8.44329×10^{-17}
0.4	3.0799×10^{-17}	1.35734×10^{-16}
0.5	1.05918×10^{-16}	3.27963×10^{-16}
0.6	9.68481×10^{-17}	4.29915×10^{-16}
0.7	1.47862×10^{-16}	4.80929×10^{-16}
0.8	2.97549×10^{-17}	4.14334×10^{-16}
0.9	2.47817×10^{-16}	2.57721×10^{-17}
1.0	6.64357×10^{-17}	6.64357×10^{-16}

$y_1(0) = 2, y'_1(0) = -12, y''_1(0) = 20$
$y_2(0) = -2, y'_1(0) = 28, y''_1(0) = -52$
$y_3(0) = -12, y'_3(0) = -33, y''_1(0) = 5$

and exact solutions

$$y_1(x) = e^x - 2e^{2x} + 3e^{-3x}$$
$$y_2(x) = 3e^x + 2e^{2x} - 7e^{-3x}$$
$$y_3(x) = -11e^x - 5e^{2x} + 4e^{-3x}$$

We compare our method m = 3 of order 5 with the BVM5 in Ref. [11] and ITPBO9 in Ref. [5] of orders 6 and 9 respectively. From the results presented in Table 17, our method m = 3 performs favourably well against the BVM5 and ITPBO9. Table 18 shows the maximum error of our methods m = 3, 5, 7. This problem is solved with various values *N*.

3.8. Problem 8

Consider the non-linear system

$$\begin{cases} y_1^{'''}(x) = \frac{1}{2}exp(4x) y_3(x) y_2'(x) \\ y_2^{'''}(x) = \frac{8}{3}exp(2x) y_1(x) y_3'(x) , \quad (82) \\ y_3^{'''}(x) = 27 y_2(x) y_1'(x) \end{cases}$$

which was solved in Ref. [5], with the initial conditions

Table 17. MAE Comparison methods $m = 3$, BVM5 and ITPBO9 (Prob			and ITPBO9 (Problem 7)
N	<i>m</i> = 3	BVM5	ITPBO9
10	1.51×10^{-7}	5.45×10^{-3}	1.42×10^{-12}
20	2.36×10^{-9}	9.59×10^{-5}	9.52×10^{-13}
40	3.68×10^{-11}	1.80×10^{-6}	5.40×10^{-13}
80	5.76×10^{-13}	2.98×10^{-8}	3.13×10^{-13}
160	8.99×10^{-15}	5.29×10^{-10}	1.22×10^{-13}
320	1.41×10^{-16}	1.07×10^{-11}	not given

Table 18. MAE methods m = 5 and m = 7 with various N (Problem 7)

Ν	m = 3	m = 5	m = 7
6	3.21×10^{-6}	1.16×10^{-9}	3.13×10^{-13}
10	1.51×10^{-7}	1.96×10^{-11}	1.90×10^{-15}
12	5.05×10^{-8}	4.57×10^{-12}	3.08×10^{-16}
20	2.36×10^{-9}	7.69×10^{-14}	1.86×10^{-18}
24	7.89×10^{-10}	1.79×10^{-14}	3.00×10^{-19}
40	3.68×10^{-11}	3.00×10^{-16}	1.82×10^{-21}
48	1.23×10^{-11}	6.99×10^{-17}	2.94×10^{-22}
80	5.76×10^{-13}	1.17×10^{-18}	1.78×10^{-24}
96	1.93×10^{-13}	2.73×10^{-19}	2.87×10^{-25}
160	8.99×10^{-15}	4.58×10^{-21}	1.74×10^{-27}
192	3.01×10^{-15}	1.07×10^{-21}	2.80×10^{-28}
320	1.41×10^{-16}	1.79×10^{-23}	1.70×10^{-30}

$$\begin{pmatrix} y_1(0) = 1, y'_1(0) = -1, y''_1(0) = 1 \\ y_2(0) = 1, y'_1(0) = -1, y''_1(0) = 1 \\ y_3(0) = 1, y'_3(0) = -3, y''_1(0) = 9 \end{pmatrix}$$

and exact solutions

$$y_1(x) = e^x$$

$$y_2(x) = e^{-2x}$$

$$y_3(x) = e^{-3x}$$

	Table 19. AE at grid points for method $m = 3$ (Problem 8)				
x	<i>y</i> 1	<i>y</i> 2	У3		
0.1	1.51546×10^{-14}	3.83364×10^{-12}	8.95258×10^{-11}		
0.2	3.10129×10^{-14}	1.73926×10^{-11}	3.39825×10^{-10}		
0.3	1.7698×10^{-13}	4.74889×10^{-11}	7.1774×10^{-10}		
0.4	1.31187×10^{-12}	1.06305×10^{-10}	1.19818×10^{-9}		
0.5	4.92628×10^{-12}	2.11419×10^{-10}	1.76148×10^{-9}		
0.6	1.3897×10^{-11}	3.85917×10^{-10}	2.39116×10^{-9}		
0.7	3.30345×10^{-11}	6.58552×10^{-10}	3.0719×10^{-9}		
0.8	6.98675×10^{-11}	1.06386×10^{-9}	3.78767×10^{-9}		
0.9	1.35634×10^{-10}	1.64216×10^{-9}	4.51995×10^{-9}		
1.0	2.46510×10^{-10}	2.43917×10^{-9}	5.2460×10^{-9}		

	Table 20. AE at grid points for method $m = 5$ (Problem 8)				
x	<i>y</i> 1	<i>y</i> 2	У3		
0.1	5.55827×10^{-17}	4.81773×10^{-16}	1.29796×10^{-14}		
0.2	2.95383×10^{-16}	1.52751×10^{-15}	4.95592×10^{-14}		
0.3	8.98209×10^{-16}	4.7062×10^{-15}	1.04197×10^{-13}		
0.4	1.91418×10^{-15}	1.1297×10^{-14}	1.72933×10^{-13}		
0.5	2.99826×10^{-15}	2.41904×10^{-14}	2.53058×10^{-13}		
0.6	5.17479×10^{-15}	4.68669×10^{-14}	3.41539×10^{-13}		
0.7	8.61645×10^{-15}	8.32541×10^{-14}	4.3678×10^{-13}		
0.8	1.46822×10^{-14}	1.38331×10^{-13}	5.36118×10^{-13}		
0.9	2.44228×10^{-14}	2.1786×10^{-13}	6.38609×10^{-13}		
1.0	4.11462×10^{-14}	3.27866×10^{-13}	7.38401×10^{-13}		

Tables 19, 20 and 21 shows the AE at selected points in the integration domain for the problem solved with step size $h = \frac{1}{10}$ and compared with the theoretical solution. The numerical results for this problem are presented in Tables 19, 20 and 21.

Table 21. AE at grid points for method m = 7 (Problem 8) х *y*₁ y_2 *y*₃ 1.48706×10^{-16} 5.55827×10^{-17} 1.21053×10^{-16} 0.1 0.2 2.95383×10^{-16} 3.06269×10^{-16} 1.54285×10^{-16} 7.87186×10^{-16} 8.20419×10^{-16} 3.35219×10^{-16} 0.3 1.35907×10^{-15} 8.6094×10^{-16} 6.26151×10^{-16} 0.4 2.44315×10^{-15} 1.26433×10^{-15} 6.49266×10^{-16} 0.5 3.39843×10^{-15} 1.79189×10^{-15} 3.95042×10^{-16} 0.6 3.1434×10^{-16} 4.67516×10^{-15} 0.7 2.37434×10^{-15} 0.8 5.63386×10^{-15} 3.02269×10^{-15} 1.57661×10^{-15} 3.53019×10^{-15} 0.9 6.93674×10^{-15} 3.58693×10^{-15} 7.89501×10^{-15} 3.95862×10^{-15} 6.27424×10^{-15} 1.0

4. Conclusion

We proposed a class of single-step hybrid block methods for solving third-order ODEs without first converting to an analogous first-order system. The convergence of the new methods was established, and the schemes are implemented without the use of starting values or predictors, avoiding the necessity for complex subroutines. To illustrate the effectiveness of the new methods, eight test problems of varied degrees of difficulty are considered. From the results, it is shown that the improvement in accuracy increases as the number of intra-step points m is increased. Tables 1-21 discuss the numerical outcomes in further depth. Our future research will focus on developing optimized single-step hybrid block schemes for solving third-order ODEs directly.

Acknowledgment

We thank the referees for the positive enlightening comments and suggestions, which have greatly helped us in making improvements to this paper. This research is supported by the Department of Science and Innovation - National Research Foundation Centre of Excellence in Mathematical and Statistical Sciences (DSI-NRF CoE-MaSS), South Africa.

References

 P. Henrici, Discrete variable methods in ordinary differential equations, Wiley, New York, 1962.

- [2] M. Rajabi, F. Ismail & N. Senu "Linear 3 and 5-step methods using Taylor series expansion for solving special 3rd order ODEs", AIP Conference Proceedings 1739 (2016) 020036. https://doi.org/10.1063/1.4952516.
- [3] O. Adeyeye & Z. Omar, "Solving third order ordinary differential equations using one-step block method with four equidistant generalized hybrid points", IAENG International Journal of Applied Mathematics 49 (2019) 1. https://www.iaeng.org/IJAM/issues_v49/issue_2/IJAM_49_2_ 15.pdf.
- [4] Z. A. Majid, N. A. Azmi, M. Suleiman & Z. B. Ibrahim, "Solving directly general third order ordinary differential equations using two-point four step block method", Sains Malaysiana 41 (2012) 623. http://www.ukm. edu.my/jsm/pdf_files/SM-PDF-41-5-2012/1520Zanariah.pdf.
- [5] R. Allogmany & F. Ismail, "Implicit Three-Point Block Numerical Algorithm for Solving Third Order Initial Value Problem Directly with Applications", Mathematics 8 (2020) 1771. https://doi.org/10.3390/ math8101771.
- [6] S. N. Jator, "On the numerical integration of third order boundary value problems by linear multistep methods", International Journal of Pure and Applied Mathematics 46 (2008) 375. https://www.researchgate.net.
- [7] K. Hussain, O. Adeyeye & N. Ahmad, "Numerical Solution of Second Order Fuzzy Ordinary Differential Equations using Two-Step Block Method with Third and Fourth Derivatives", Journal of the Nigerian Society of Physical Sciences 5 (2023) 1087. https://doi.org/10.46481/jnsps.2023. 1087.
- [8] O. Adeyeye, A. Aldalbahi, J. Raza, Z. Omar, M. Rahaman, M. Rahimi-Gorji & N. M. Hoang, "Nonlinear solution of the reaction-diffusion equation using a two-step third-fourth derivative block method", International Journal of Nonlinear Sciences and Numerical Simulation 22 (2021) 111. https://doi.org/10.1515/ijnsns-2019-0309.
- [9] D. O. Awoyemi & O. M. Idowu, "A class of hybrid collocation methods for third-order ordinary differential equation", International Journal of Computer Mathematics 82 (2005) 1287. https://doi.org/10.1080/ 00207160500112902.
- [10] J. Sunday & J. N. Ndam, "An Accuracy-preserving Block Hybrid Algorithm for the Integration of Second-order Physical Systems with Oscillatory Solutions", Journal of the Nigerian Society of Physical Sciences 5 (2023) 1017. https://doi.org/10.46481/jnsps.2023.1017.
- [11] S. N. Jator, T. Okunlola & T. Biala, R. Adeniyi, "Direct integrators for the general third-Order ordinary differential equations with an application to the Korteweg–de Vries Equation", International Journal of Applied & Computational Mathematics 4 (2018) 1. https://doi.org/10.1007/ s40819-018-0542-6.
- [12] S. N. Jator & E. Agyingi, "Block Hybrid-Step Backward Differentiation Formulas for Large Stiff Systems", International Journal of Computational Mathematics 2014 (2014) 8. https://doi.org/10.1155/2014/162103.
- [13] H. Ramos, "An optimized two-step hybrid block method for solving first-order initial value problems in ODEs", Differential Geometry- Dynamical System 19 (2017) 107. http://vectron.mathem.pub.ro/dgds/v19/ D19-ra-C97P.pdf.
- [14] M. G. Orakwelu, S. Goqo & S. Motsa, "An optimized two-step block hybrid method with symmetric intra-step points for second order initial value problems", Engineering Letters 29 (2021) 948. http://www. engineeringletters.com/issues_v29/issue_3/EL_29_3_18.pdf.
- [15] U. Mohammed, J. Garba, A. A. Alhaji, A. I. Maáli & M. Semenov, "Second derivatives single step block hybrid method for nonlinear dynamical systems", Nigerian Journal of Mathematics and Applications **32** (2022) 42. http://repository.futminna.edu.ng:8080/jspui/handle/ 123456789/19414.
- [16] P. Onumanyi, U. W. Sirisena & S. N. Jator, "Continuous finite difference approximation for solving differential equations", International Journal of Computer Mathematics 72 (1999) 15. https://doi.org/10.1080/ 00207169908804831.
- [17] S. N. Jator, "On a class of hybrid method for y'' = f(x, y, y')", International Journal of Pure and Applied Mathematics **59** (2010) 381. https://doi.org/10.1080/00207169108804026.